

Group and representation theory: A Shitty summary

Chapter 2: Groups: Basic definitions and properties

2.1) Basics: finite groups

def: A group is a set of elements which can transform into each other associatively $[(a \cdot b) \cdot c = a \cdot (b \cdot c) \dots]$ via a group action (not the GAI)

$\phi: G \times G \rightarrow G$. A group must contain: an identity element, and every element must have an inverse: $\forall g \in G, \exists g^{-1} \in G: g \cdot g^{-1} = g^{-1} \cdot g = 1$

def: we can define 2 types of group actions

let G be a group, X be some finite set
left group action:

$d: G \times X \rightarrow X$ if d satisfies: $\forall x \in X: d(e, x) = x$ identity

~~$\forall g, h, x \in X: d(g, d(h, x)) = d(gh, x)$~~
 $\forall g, h \in G, x \in X: d(d(x, g), h) = d(x, gh)$

we say the group acts on X (from the left).

A set X combined with some action of G is called a (left) G -set.

Right group action:

$r: X \times G \rightarrow X$ if r satisfies:

$\forall x \in X: r(x, e) = x$

$\forall g, h \in G, x \in X: r(r(x, g), h) = r(x, gh)$

If the action is commutative then $L\Lambda = R\Lambda$.

We can say the action of a group G on some set X is:

transitive: $\forall x, y \in X, \exists g \in G: g \cdot x = y$

faithful: $\forall g, h \in G, \exists x \in X: g \cdot x \neq h \cdot x$

free: $\forall g, h \in G, x \in X: g \cdot x = h \cdot x \Leftrightarrow g = h$ (free action is always faithful)

regular: transitive + free

:

the number of elements in a group is called it's cardinality. (denoted $|G|$ or $\#G$)

We can display the way 2 group elements combine through their action with a Cayley table.

Let $\{a, b, c\}$ be some group, it's Cayley table is then

	a	b	c
a	a^2	ab	ac
b	ba	b^2	bc
c	ca	cb	c^2

We can derive some interesting properties from this table:

A Abelian group \Leftrightarrow Cayley table is symmetric along its diagonal

A very important family of groups is the permutation group S_n of n elements. It has cardinality $n!$ and is also called the symmetric group.

We can also choose to define our groups through a generating set of group elements, that through multiplication, result in all the elements of that group

Example: the dihedral group of order n is generated by (τ, σ) where

τ is a rotation over 180° and σ is any reflection about a line of symmetry.

The upper limit on the size of this set is $\log_2(|G|)$, where this is only satisfied by a generating set created by direct products of \mathbb{Z}_2 's

2.1.2) Automorphisms

Def:

An automorphism of a group is a permutation (read: relabeling) of its elements $a: G \rightarrow G$ s.t. the Cayley table remains invariant under that relabeling

example:

the regular Cayley table of the Klein four-group (every element is its own inverse):

$$\begin{array}{c} e \ a \ b \ c \\ a \ e \ c \ b \\ b \ c \ a \ e \\ c \ b \ a \ e \end{array} \xrightarrow{f} \begin{array}{c} e \ c \ a \ b \\ c \ e \ b \ a \\ a \ b \ e \ c \\ b \ a \ c \ e \end{array}$$

and the second is under some automorphism f given by: $f(e) = e$

$$f(a) = c$$

$$f(b) = a$$

$$f(c) = b$$

another way to define the automorphism is by saying:

$$f \in \text{Aut}(G) \quad \forall g_1, g_2 \in G : f(g_1) f(g_2) = f(g_1 g_2)$$

the automorphisms of a group G themselves also form a group

2.1.3) Classes and Subgroups

def:

A subgroup H of some group G is a subset of elements of G that themselves form a group.

for some subgroup H of some group G we define the left coset $aH = \{ah \mid h \in H\}$ for some $a \in G$. the amount of left cosets is denoted by $[G:H]$ and is called the index. (analogous for right cosets).

Theorem:

Lagrange's theorem states that for some subgroup H of G :

$$[G:H] = \frac{|G|}{|H|}$$

(thus that the cardinality of H is a divisor of the cardinality of G)

Proof:

the left cosets are equivalence classes of some relation on G . call x and y in G equivalent if $\exists h \in H \ni x = hy$. the left cosets thus define a partition of G . each left coset aH must have the same cardinality as H because $x \mapsto ax$ defines a bijection $H \xrightarrow{\sim} aH$. thus

$$|G| = [G:H] \cdot |H|$$

def:

the centre of some group G is the subgroup of elements which commute with all other elements of G . (Also called the centraliser, denoted $Z(G)$)
(trivially found by looking for zero rows in $C - C^T$ where C is the entry table.)

another way of constructing interesting subgroups is by creating classes through conjugation: for some g we create the set $\{hgh^{-1} \mid h \in G\}$. we can divide the group in such classes and every element is contained in exactly 1 class.

if we now create a subgroup of complete classes this defines a group closed under conjugation, which we call a normal/invariant subgroup (note the centre is a normal subgroup)

a subgroup N is normal if thus it is invariant under conjugation with all group elements of its 'parent' group. we also have $\forall g \in G : gN = Ng$.

this gives rise to some types of groups:

- simple :

a group is called simple if it's only normal subgroup is the set consisting of solely the identity $\{e\}$

- factor :

the factor (or quotient) group G/N is the group of cosets aN

- maximal normal subgroup:

a normal subgroup that contains no non-trivial subgroups
(it's factor group is thus simple)

- commutator group:

this is the group generated from all group elements $xgx^{-1}g^{-1}$. this allows us to define the Abelianization $G/[G, G]$.

2.7.4) extending groups

Semi-direct product

for two groups N and H the semi-direct product of N and H w.r.t ρ

$$G = N \rtimes_{\rho} H$$

is the group of elements $N \times H$ i.e. all elements of G are as (n, h) with group multiplication defined as

$$(n_1, h_1) \cdot (n_2, h_2) := (n_1 \rho_{h_1}(n_2), h_1 h_2) \quad (\rho_{h_1} = \rho(h_1))$$

with ρ a homomorphism $\rho: H \rightarrow \text{Aut}(N)$

then there is a normal subgroup of G (H not necessarily, but it is always a subgroup)
(note that this only defines G up to isomorphism)

Sequences and extensions (go ahead and skip this)

def:

an exact sequence is a sequence of homomorphisms where the image of one morphism is the kernel of the next. i.e.

$$\text{im}(\phi_i) = \ker(\phi_{i+1}) \quad \forall i$$

$$(\phi_i: G_i \rightarrow G_{i+1})$$

def:

a short exact sequence is an exact sequence of the form

$$1 \rightarrow A \xrightarrow{f} B \xrightarrow{g} B/A \rightarrow 1$$

a special type of short sequence is called a split short sequence.

this is when $\exists h: B/A \rightarrow B : g \circ h = \text{id}_{B/A}$

A good question to ask now is, for some given N and H find all G compatible with

$$1 \rightarrow N \rightarrow G \rightarrow H \rightarrow 1$$

as it turns out this is a very difficult problem. But if we demand that this is a split short sequence, then all constructions are found through

the semi direct product:
 $N \rtimes H$

we call these the splitting extensions

another type of interesting extension is the central extension:

def:

central extensions G of some group H by an Abelian group A are defined by:

$$1 \rightarrow A \rightarrow G \rightarrow H \rightarrow 1$$

where we require A to be a subgroup of the centre $Z(G)$ and where
 $H = G/A$

all central extensions can be found by defining elements (a, h) and
a group multiplication rule:

$$(a_1, h_1) \cdot (a_2, h_2) = (a_1 a_2 + w(h_1, h_2), h_1 h_2).$$

Rest of chapter 2 talks about cohomology. This is not expected to be known.

Chapter 3: Representation theory.

up till now i have been rewriting / clarifying the course notes doc.
seeing as how go% of it isn't expected to be known i shall now lose myself off of the theory included in the exercise session documents.

3.1) definition and left/right regular repr.

def:

a representation of some finite group G is a group homomorphism

$\varphi: G \rightarrow GL(V)$, where V is a finite vector space. i.e. φ maps every group element g to some matrix X_g s.t.

$$X_{g_1} X_{g_2} = X_{(g_1 g_2)} \quad \forall g_1, g_2 \in G$$

we define the dimension of the representation as $\dim(V)$

we call a representation faithful if S is injective (every element has at least 1 unique value)

And now, back to the lecture notes :

note that $S(e) = \mathbb{1}$. This implies that all these matrices X_g must be invertible!
even stronger still: for finite groups we will always be able to restrict ourselves to unitary matrices ($A^{-1} = A^*$)

Finally, it holds that for any 2 equivalent representations (that is to say, 2 repr. related through conjugation: $A_g = X(g)X^{-1}$) they are related by a unitary conjugation.

Now to construct these representations.

one clear way is by using the regular representation, which is obtained by constructing $n \times n$ permutation matrices which encode the information of the Cayley table (thus how the elements transform into each other) of the group with cardinality n .

Left regular:

the left regular representation is given by

$$L_g = \sum_h |g \cdot h\rangle \langle h|$$

$|g\rangle$: group element

$$\text{for which } L_x \cdot L_y = L_{xy}.$$

$|g \cdot h\rangle \langle h|$ is $h \mapsto g \cdot h$ end does not mix ordering elem

Right regular:

this is given by

$$R_g = \sum_h |h \cdot g\rangle \langle h|$$

$$\text{for which } R_x R_y = R_{y \cdot x}$$

3.0.1) Irreducible representations: Basics

def: invariant subspace:

A subspace Y is called invariant if $\forall A_g$ in the algebra

$$A_g Y \subset Y \quad (\text{left-invariant})$$

$$YA_g \subset Y \quad (\text{right-invariant})$$

for matrix-algebras with invariant subspaces there always exists a basis on which we can write these matrices as block diagonal

$$A(z) = \begin{bmatrix} B(z) & D(z) \\ 0 & C(z) \end{bmatrix} \text{ with } \begin{bmatrix} X \\ 0 \end{bmatrix} \quad (\text{left-inv})$$

$$A(z) = \begin{bmatrix} B(z) & 0 \\ D(z) & C(z) \end{bmatrix} \quad (\text{right-inv})$$

there now exists a theorem (Mazurkevich's theorem) that says:

"There always exists a similarity transform ($P^{-1}AP$) which makes all the matrices of the representation block diagonal"

Proof:

$$\text{starting from: } A(z) = \begin{pmatrix} B(z) & 0 \\ D(z) & C(z) \end{pmatrix}$$

we want:

$$\begin{pmatrix} I & 0 \\ -F & I \end{pmatrix} \begin{pmatrix} B(z) & 0 \\ D(z) & C(z) \end{pmatrix} \begin{pmatrix} I & 0 \\ F & 0 \end{pmatrix}$$

to be diagonal.

$$\text{this is equal to: } \begin{pmatrix} B(z) & 0 \\ -FD(z) + D(z) + C(z)F & C(z) \end{pmatrix}$$

or simply:

$$-FD(z) + D(z) + C(z)F = 0$$

$$\text{let } F = \frac{1}{|G|} \sum_{a \in G} D(a) B(a^{-1})$$

$$C(z)$$

then:

$$-\frac{1}{|G|} \sum_{a \in G} D(a) B(a^{-1}) B(z) + D(z) + \frac{1}{|G|} \sum_{a \in G} D(a) B(a^{-1})$$

but

$$D(z)a = D(z)B(a) + C(z)D(a) \leftrightarrow C(z)D(a) = D(z)a - D(z)B(a)$$

substituting this we find

$$\begin{aligned}
 0 &\stackrel{?}{=} -\frac{1}{|G|} \sum D(a) D(a^{-1}) B(a) + D(e) + \frac{1}{|G|} \sum (D(a) - D(e) B(a)) B(a^{-1}) \\
 &= -\frac{1}{|G|} \underbrace{\sum_{a \in G} D(a) B(a^{-1}) D(a)}_{= D(e)} + D(e) + \frac{1}{|G|} \sum D(a) B(a^{-1}) - \underbrace{\frac{1}{|G|} \sum_{a \in G} D(a) B(a) B(a^{-1})}_{= D(e)} \\
 &= -\frac{1}{|G|} \sum_{a \in G} D(a) B(a^{-1}) B(e) + \frac{1}{|G|} \sum_{a \in G} D(a) B(a^{-1})
 \end{aligned}$$

where the two sums are equal because of the rearrangement theorem
and thus we have proven it \square

thus any representation $U(g)$ can be written in which

$$U(g) = \begin{pmatrix} U_1(g) & 0 & \cdots & 0 \\ 0 & U_2(g) & & \vdots \\ \vdots & 0 & \ddots & \vdots \\ 0 & \vdots & \vdots & \vdots \end{pmatrix}$$

where the $U_i(g)$'s don't have invariant subspaces!

Def: irreducible representation (irrep):

an irrep is a matrix representation of a group with no invariant subspaces.

a crucial tool will be Schur's lemma:

irrep spans matrix-alg. \rightarrow 1: if $U(g)$ is an irrep and $\forall g : X U(g) = U(g) X \Rightarrow X = c \cdot I$

2: given 2 irreps $U(g), V(g)$ and a matrix $T \in \mathbb{C}^{|G| \times |G|}$ $\forall g : T U(g) = V(g) T$ then U, V equivalent or $T = 0$

for finite groups, any irrep is always equivalent to an irrep of unitary matrices

Proof:

see lecture notes (is not expected to be known)

a second point is that representations can be written as a direct sum of irreps
this as a direct consequence of maschke's theorem. block diagonalisable \rightarrow \oplus irreps

The great orthogonality theorem

Consider all N inequivalent unitary irreps of G : $D_{ij}^{\alpha}(x)$ (d_{α} = matrix dim.)
(α labels the irreps). Then:

$$\frac{1}{|G|} \sum_i D_{ij}^{\alpha}(x) \bar{D}_{kl}^{\beta}(x) = \frac{1}{d_{\alpha}} \delta_{jk} \delta_{il} \delta_{\alpha \beta}$$

(Proof not expected to be known)

This implies:

$$\sum_{\alpha} d_{\alpha}^2 \leq |G|$$

Proof:

Let $i=j, l=k$ in G.O.T

$$\frac{1}{|G|} \sum_i D_{ii}^{\alpha}(x) \bar{D}_{kk}^{\beta}(x) = \frac{1}{d_{\alpha}} \delta_{\alpha \beta} \delta_{kk}$$

rewritten:

$$\frac{1}{|G|} \sum_{i \in \alpha} \sum_{k \in \beta} D_{ik}^{\alpha} \bar{D}_{ik}^{\beta} = \frac{1}{d_{\alpha}} \delta_{\alpha \beta} \sum_{i \in \alpha} \sum_{k \in \beta} \delta_{ik} = \delta_{\alpha \beta}$$

Sums over i and k are the traces of D^{α} and \bar{D}^{β} , which we define as the characters of these representations

$$\frac{1}{|G|} \sum_{i \in \alpha} \sum_{k \in \beta} D_{ik}^{\alpha} \bar{D}_{ik}^{\beta} = \frac{1}{|G|} \sum_{i \in \alpha} \overline{\text{Tr}(D^{\alpha}(i))} \text{Tr}(\bar{D}^{\beta}(i)) = \underbrace{\frac{1}{|G|} \sum_{i \in \alpha} \chi_{\alpha}(i) \bar{\chi}_{\beta}(i)}_{\text{G.O.T for characters}} = \delta_{\alpha \beta}$$

characters of representations

where irrs are rather arbitrary (defined up to equivalence relations)

char traces are basis invariant, and therefore universal we thus define

$$\boxed{\chi^{\alpha}(x) = \text{Tr}(D^{\alpha}(x))}$$

the group character is constant in a given class (i.e conjugation with other group elements) Let us denote the classes as C_k and define N_k as the number of elements in that class

from before we know, the G.O.T for characters:

$$\frac{1}{|G|} \sum_{\lambda} \chi^{(\alpha)}(\lambda) \chi^{(\beta)}(\lambda) = \delta_{\alpha \beta}$$

they thus form an orthonormal basis in the space of group elements

we can construct a matrix with orthonormal columns O_{xa} :

$$\sqrt{\frac{N_k}{|G|}} \chi^a(x_k)$$

which shows:

- $|C_k| \geq |G|$ (otherwise O_M not full rank)
- if 2 chars have the same characters \Rightarrow they are equivalent

we can thus readily calculate how many times a reducible rep. contains an irrep: calculate the inner prod of the corresponding characters

in particular for the left-regular repr. say an irrep α :

$$C_\alpha = \frac{1}{|G|} \sum_x \chi^\alpha(x) \bar{\chi}^\alpha(x) = d_\alpha \quad (L(g) = \sum_h g \cdot h \chi^\alpha(h))$$

(d_α = dim of irrep α)

note $\bar{\chi}^\alpha(1) = d_\alpha$. This now implies that the left-regular representation contains each d_α dim-irrep $n_\alpha = d_\alpha$ times!

this now implies

$$\sum_\alpha d_\alpha^2 = |G|$$

and that $\sum C_\alpha = |G|$

see next for example.

Real, complex and quaternionic representations, the schur indicator and kramer's theorem

we can classify irreps into 3 categories. By defining an anti-unitary operator J for which: $J(\sum a_i x_i) = \sum \bar{a}_i J(x_i)$. we define the complex conjugate repr.

\bar{U} of U as:

$$\bar{U}(g) = J U(g) J \quad (J^2 \text{ is } I \Rightarrow \text{is also a repr})$$

a repr is self conjugate if \bar{U} is equivalent to U . i.e.: $\exists X: \bar{U} = X U X^{-1}$

if there exists no such $X \rightarrow$ repr is complex

taking the conjugate of

$$\bar{U} = X U X^{-1}$$

gives

$$U(g) = \bar{X} X U(g) (\bar{X} X)^{-1} \quad (1)$$

schur's lemma says $\bar{X} X = C \cdot I$ and $X = C X^T$

taking the conjugate of (1) then

$$U(g) = C^2 \bar{U}(g)$$

remember U to be self conjugate $\Rightarrow C^2 = 1 \Rightarrow C = \pm 1$ which distinguishes the 2 other classes:

real, $C=+1: X=X^T$, thus with some decomposition we write

$$X = QQT^T. \text{ we want some } Y \text{ s.t.: } \overline{Y U(g) Y^{-1}} = Y U(g) Y^{-1}$$

$$\text{or } Y X U(g) X^{-1} Y^{-1} = Y U(g) Y^{-1}. Y = Q^T \text{ is just that.}$$

Proof: $C=-1:$

wont be asked, wont write it down

a simple way to verify what type of repr you have is through its characters

$$\frac{1}{|G|} \sum_{x \in G} \chi(x^2) = \begin{cases} 1 & \text{real} \\ 0 & \text{complex} \\ -1 & \text{quot} \end{cases}$$

Proof:

$$\frac{1}{|G|} \sum_{x \in G} \chi(x^2) = \frac{1}{|G|} \sum_{g \in G} \sum_{y \in g} \chi(y) \overline{\chi(y)} = \frac{1}{|G|} \sum_{g \in G} \sum_{y \in g} \chi(y) \bar{\chi}(y)$$

if U is real the sum reduces to $\frac{1}{d} \sum_{g \in G} d_g d_g = d$

if U is complex, the repr are inequivalent $\Rightarrow \sum_{g \in G} d_g d_g = 0$

if U is quot, $U = X^T \bar{U} X, \bar{U} X = -I$, sum becomes

$$\frac{1}{|G|} \sum_{g \in G} \text{tr}(X^T \bar{U}(g) X \bar{\chi}(g)) = \frac{1}{d} \sum_{g \in G} \bar{X} X_{gg} = \frac{1}{d} \text{tr}(X^T X) = -1$$

Projective representations

The central idea here is that if $U(g) \otimes U(g)$ forms a representation, $U(g)$ does not always form a repr. if they satisfy

$$U(g)U(h) = e^{iw(gh)} U(gh) \quad w(g, h) \text{ some real phase}$$

then $U \otimes U$ is a repr but U isn't, we call U a projective repr.

A general projective repr is defined as one where its product is not compatible with the group action but results in an extra prefactor

$$\rho(g)\rho(h) = f(g, h) \rho(gh)$$

We impose the condition on w that:

$$w(x, g) + w(x, y, z) \underset{\text{mod } 2\pi}{=} w(x, y, z) + w(y, z) \quad (2) \quad (\text{mod } 2\pi \text{ means equivalent in mod } 2\pi)$$

This implies there are only a finite amount of projective repr!

We will derive (2):

$$U(x)U(g)U(z) = U(x)(e^{iw(y,z)}U(yz)) = e^{iw(x,yz)+iw(y,z)}U(xyz)$$

or representation is always associative (odds):

$$U(x)[U(y)U(z)] = [U(x)U(y)]U(z) = e^{iw(x,y)+iw(xy,z)}U(xyz)$$

thus not

$$w(x, yz) + w(y, z) \underset{\text{mod } 2\pi}{=} w(x, y) + w(xy, z) \quad (e^{cx} = e^{c(x+2\pi)})$$

Constructing Projective repr.

For some given cocycle $w(g, h)$ we find the corresponding projective repr by starting from the regular representation

$$R_w = \sum_h e^{iw(g,h)} |g \cdot h\rangle \langle h|$$

indeed:

$$\begin{aligned} R_w(g)R_w(h) &= \sum_{x,y} e^{(\omega(y,x) + \omega(h \cdot y))} |g \cdot x\rangle \langle h \cdot y| \\ &= \sum_y e^{(\omega(y, h \cdot y) + \omega(h, y))} |g \cdot h \cdot y\rangle \langle y| \\ &= \sum_y e^{(\omega(g \cdot h, y) + \omega(g, h))} |g \cdot h \cdot y\rangle \langle y| \\ &= e^{\omega(g, h)} R_w(g \cdot h) \end{aligned}$$

$\Rightarrow x = hy$, order ∞ (orthonormal)

Note, projective representations with non-trivial cocycles can't have 1 dim subspaces. Otherwise, let IV be this inv. subspace:

$$R_w(g) \stackrel{IV}{=} e^{\omega(g)} IV$$

$$\begin{aligned} R_w(g)R_w(h) &= e^{(\omega(g) + \omega(h))} IV \\ &= e^{\omega(g, h)} R(g \cdot h) IV \\ &= e^{\omega(g, h)} e^{\omega(g)} (R_w(g \cdot h)) IV \end{aligned}$$

↳ die volgt niet uit ell., ergens ook gelijk aan $R_w(g)R_w(h)$ uit def
Proj. repr. !!

now $\omega(g, h) = \alpha(g) + \alpha(h) - \alpha(gh)$ which is a trivial cocycle and thus a contradiction.

3.0.2: The group algebra

3.0.3: group tolleaux

3.0.4: induced rep

} no theory to be known here, if needed for some exercise

i will write it here

Chapter 9: Lie groups and their representations

9.1: Associativity

we would like to generalize now to continuous groups.

we do this by looking at group elements in a certain domain and assume that the group structure is encoded in an analytical function which satisfies associativity

$$\phi(a, \phi(b, c)) = \phi(\phi(a, b), c)$$

9.2: Lie groups and structure factors

we require the associativity mentioned above, as well as unity elements and the existence of inverses.

to develop the theory of lie groups we define a vector of variables x on which the elements of the group will act. The point of a lie group then is that its properties are completely controlled by behaviour under infinitesimal transformations

$$x' = f(x, a)$$

$$\begin{aligned} x' + da' &= f(x', da) \\ &= f(x, a + da) \end{aligned}$$

clearly

$$da_i = \underbrace{\frac{\partial \phi_i(a, b)}{\partial b_j}}_{\Phi_y(a)} \Big|_{b=0} da_j$$

we can invert this to get $\Phi(a)$

$$da = \Phi(a) da \rightarrow U_R(x')$$

also

$$dx'_i = \sum_k \frac{\partial f_i(x', a)}{\partial a_k} \Big|_{a=0} da_k$$

$$dx_i' = U(x') \Psi(u) da$$

$$\frac{dx_i'}{\delta a_k} = \int U_{ik}(x') \Psi_{jk}(u)$$

looking now at some infinitesimal transformation

$$\begin{aligned} x_i + dx_i &= f_i(x, \delta a) \\ &= x_i + \sum_R \frac{\partial f_i}{\partial a_R} \delta a_R \\ &= x_i + \sum_k U_{ik} \delta a_k \\ &= \underbrace{\left(1 + \sum_R \sum_i U_{ik} \frac{\partial}{\partial x_i} \delta a_R \right)}_{\hat{x}_k} x_i \\ &= \left(1 + \sum_R \hat{x}_R \delta a_R \right) x_i \end{aligned}$$

where we defined the generator \hat{x}_k :

$$\boxed{\hat{x}_k = \sum_i U_{ik}(x) \frac{\partial}{\partial x_i}, \quad U_{ik}(x) = \frac{\partial f_i(x, a)}{\partial a_k} \Big|_{a=0}}$$

we will study Lie groups through these generators and their commutator algebra

example 1:

$$x' = ax + b$$

$$\hat{x}_a = \frac{\partial(ax+b)}{\partial a} \cdot \frac{\partial}{\partial a} = x \frac{\partial}{\partial x}, \quad \hat{x}_b = \frac{\partial(ax+b)}{\partial b} \cdot \frac{\partial}{\partial b} = 1 \frac{\partial}{\partial x}$$

$$[\hat{x}_a, \hat{x}_b] = x \cdot \frac{\partial}{\partial x} (1 \frac{\partial}{\partial x}) - 1 \frac{\partial}{\partial x} (x \frac{\partial}{\partial x}) = 0 - \frac{\partial}{\partial x} = -\hat{x}_c$$

example 2:

$$\begin{pmatrix} x_1' \\ x_2' \end{pmatrix} = \begin{pmatrix} a & b \\ c & d \end{pmatrix} \begin{pmatrix} x_1 \\ x_2 \end{pmatrix}$$

$$\begin{cases} f_1 = ax_1 + bx_2 = x_1' \\ f_2 = cx_1 + dx_2 = x_2' \end{cases}$$

$$\hat{x}_a = \frac{\partial f_1}{\partial a} \frac{\partial}{\partial x_1} + \frac{\partial f_2}{\partial a} \frac{\partial}{\partial x_2} = x_1 \frac{\partial}{\partial x_1}, \quad \hat{x}_c = \frac{\partial f_1}{\partial c} \cdot \frac{\partial}{\partial x_1} + \frac{\partial f_2}{\partial c} \cdot \frac{\partial}{\partial x_2} = x_2 \frac{\partial}{\partial x_2}$$

$$\hat{x}_b = \frac{\partial f_1}{\partial b} \frac{\partial}{\partial x_1} + \frac{\partial f_2}{\partial b} \cdot \frac{\partial}{\partial x_2} = x_2 \frac{\partial}{\partial x_1}, \quad \hat{x}_d = \frac{\partial f_1}{\partial d} \cdot \frac{\partial}{\partial x_1} + \frac{\partial f_2}{\partial d} \cdot \frac{\partial}{\partial x_2} = x_2 \frac{\partial}{\partial x_2}$$

also see the final example on p 47.

9.2. 1) Structure factors

to derive these, let us start from

$$\frac{\partial x_i}{\partial a_j} = \sum_l U_{ij}(x) \psi_{jl}(a) \quad (3)$$

where $x_i + dx_i = f_i(x, da)$.

all these functions are analytic, s.t.:

$$\frac{\partial^2 x_i}{\partial a_k \partial a_\nu} = \frac{\partial^2 x_i}{\partial x_j \partial a_k}$$

note:

$$\frac{\partial U_{ij}(x)}{\partial a_k} = \frac{\partial U_{ij}(x)}{\partial x_j} \frac{\partial x_j}{\partial a_k} \stackrel{(3)}{=} \frac{\partial U_{ij}(x)}{\partial x_j} U_{j\nu}(x) \psi_{\nu k} \quad (\text{const. sum})$$

thus

$$U_{ij}(x) \left(\frac{\partial \psi_{\nu k}}{\partial a_\nu} - \frac{\partial \psi_{\nu l}}{\partial a_k} \right) = - \left(U_{j\nu} \frac{\partial U_{ij}}{\partial x_\nu} - U_{j\nu} \frac{\partial U_{il}}{\partial x_\nu} \right) \psi_{\nu k} \psi_{\nu l}$$

or by using the inverse

$$U_{j\nu}(x) \frac{\partial U_{ij}(x)}{\partial x_\nu} - U_{j\nu} \frac{\partial U_{ij}}{\partial x_\nu} = \underbrace{\left(\frac{\partial \psi_{\nu k}}{\partial a_\nu} - \frac{\partial \psi_{\nu l}}{\partial a_k} \right) \theta_{\nu k} \theta_{\nu l}}_{f_{ij}^k(a)} U_{ij}(x)$$

partial derivatives yield:

$$x_i: \quad \frac{\partial}{\partial x_m} \left(U_{j\nu}(x) \frac{\partial U_{ij}(x)}{\partial x_\nu} - U_{j\nu} \frac{\partial U_{ij}}{\partial x_\nu} \right) = f_{ij}^k \frac{\partial}{\partial x_m} U_{ik}(x)$$

$$a_j: \quad \frac{\partial f_{ij}^k(a)}{\partial a_g} U_{ik} = 0$$

U is full rank s.t. f_{ij}^k has no dependence on the group element and is a constant! they control the commutation relationships

$$\boxed{[\hat{x}_j(x), \hat{x}_\sigma(x)] = \sum_h f_{j\sigma}^h \hat{x}_h(x)}$$

note following properties:

- the structure factors must be anti-symmetric
- they must obey the Jacobi identity

9.2.2: adjoint representation and exponential map

describing all the groups has reduced to finding solutions $f_{\alpha\beta}^Y$ that satisfy the Jacobi identity

the representation theory consists then of finding all groups X_α satisfying the Lie commutator

$$[X_\alpha, X_\beta] = \sum_Y f_{\alpha\beta}^Y X_Y$$

for given $f_{\alpha\beta}^Y$.

groups have generators, turns out we can construct the set of generators from the structure factors. we call this solution the adjoint repr.

given an anti-symm $f_{ab}^c = -f_{ba}^c$ solution to the Jacobi eq., the adj repr is defined as:

$$\hat{X}_g := (X_g)_Y = -f_{gi}^Y$$

it's matrix elements thus are the structure factors!

9.3: classifying Lie algebras

we restrict ourselves to semi-simple, compact Lie algebras and groups, where:

semi-simple: direct sum of simple

simple body: trivial set of operators closed under commutation

compact group: allows us to take the generators of the algebra as anti-symm

9.3.7) Cartan subalgebra and Weyl basis

given the adjoint representation of some Lie algebra, is it always possible to bring the N generators into some canonical form through linear combinations.

The Cartan subalgebra is a maximal commutating algebra and its dimensionality r will be important.

Given a set of r commuting operators H_i , we can simultaneously diag. them to obtain r diagonal matrices D_i . Their diagonal elements fully characterise the Lie algebra. They have r zero entries and $N-r$ non-zero (complex) entries. We can thus define $N-r$ vectors with r entries namely $\vec{\alpha}(i)$ the i th part of $D_i(\alpha, \alpha)$. These vectors are the root-vectors rotated as $\vec{\alpha}$.

These $\vec{\alpha}$ allow us to write the algebra in the canonical Cartan-Weyl form

$$[H_i, H_j] = 0$$

$$[H_i, E_\alpha] = \alpha(i) E_\alpha$$

$$[E_\alpha, E_\beta] = c_{\alpha\beta} E_{\alpha+\beta}$$

$$[E_\alpha, E_\beta] = \sum_i \alpha(i) H_i$$

$E_{\alpha/\beta}$ = raising/lowering operators

$c_{\alpha\beta} = 0$ if $\alpha+\beta$ is not a root, some integer otherwise

these root vectors obey following symmetries:

- if α_a is a root, so is $-\alpha_a \Rightarrow E_\alpha$ can't be hermitian
- $[E_\alpha, E_\beta]$ is either an eigenvector with root $\alpha+\beta$, or zero
- if we let $\rho = -\alpha$, $[H_i, [E_\alpha, E_\beta]] = 0 \Rightarrow [E_\alpha, E_{-\alpha}]$ is in the Cartan subalgebra

$$[E_\alpha, E_\beta] = \sum_i c_i H_i$$

if H_i orthonorm:

$$c_i := \alpha(i)$$

a crucial insight now, $\forall \vec{\alpha}$ there exists a set of 3 operators

in the complete algebra, which form a subalgebra of $SU(2)$

this means the generators E_α fall into different sets which transform according to an irrep of $SU(2)$ under the adjoint of the subalgebra

$\mathcal{D}_{\pm}^{\alpha} :$

$$\mathcal{D}_2^{\alpha} = \frac{1}{\| \alpha \|_2^2} \sum_{i=1}^r \alpha(i) H_i$$

$$\mathcal{D}_{\pm}^{\alpha} = \frac{1}{\| \alpha \|_2} E_{\pm \alpha}$$

- a vector root has no multiplicities
- if $\alpha(i)$ is a root, $\lambda \alpha(i)$ is only a root for $\lambda = -1$

Highest weight and numerology of roots

we shall now construct an irrep. we pick a largest eigenvector (highest weight) w of the cartan subalgebra and then apply lowering operators to construct the irrep. we say $\vec{\alpha} > \vec{\beta}$ if the first non-zero component of $\alpha - \beta$ is > 0

all weights $w(i)$ are defined as eigenvalues of H_i :

$$H_i |w\rangle = w(i) |w\rangle$$

starting from some $w(i)$ we apply the raising and lowering operators to obtain the eigenvals $w(i) \in \mathbb{R}$ with corresp. eigenvectors $|h\rangle = E_{\pm \alpha}^{\vec{\alpha}} |w\rangle$

these $|h\rangle$ now have to span an irrep of $SU(2)$

we find

$$\frac{\vec{w} \cdot \vec{\alpha}}{\|\alpha\|_2^2} = -\frac{n}{2} \quad \text{with } n \text{ some integer}$$

note the severe restriction on angles of root vectors!

this also implies Weyl-reflection symmetry

$$\vec{w} \rightarrow w - 2 \frac{\vec{w} \cdot \vec{\alpha}}{\|\alpha\|_2^2} \vec{\alpha}$$

we could have done this with the adjoint rep. the eigenvectors then are the generators themselves

and $\omega(i) = \alpha(i)$. Thus for any 2 roots

$$\frac{\vec{P}\vec{\alpha}}{\|\alpha\|^2} = -\frac{n'}{2} \quad \frac{\vec{\alpha}\vec{P}}{\|\alpha\|^2} = -\frac{n''}{2}$$

which greatly restricts all possible roots as:

$$\frac{(\vec{P}\vec{\alpha})^2}{\|\alpha\|^2 \|\vec{P}\|^2} = \cos^2(\theta_{\alpha P}) \quad \text{and} \quad \frac{\|\alpha\|^2}{\|\vec{P}\|^2} = \frac{n''}{n'}$$

which results

n'	θ	α^2/\vec{P}^2
0	90	any
1	60/120	1
2	45/135	1/2, 2
3	30, 150	1/3, 3
4	0	1

all that remains is to find vectors which satisfy these conditions

See pg 9 for the example in $SU(3)$ (also see the root diagrams)

dynkin diagrams

theory questions
end here :)

there is redundancy in root diagrams, we can always transform most of them into each other through linear combination. there exists a systematic way of selecting a subset of vectors that characterize all roots we call them the simple vectors (simple roots).

they are by definition positive and the size of this set is the dimension of the Cartan subalgebra: r

the only possible angles between these is: $90, 120, 135, 150$.

we follow the following rules when drawing a dynkin diagram of simple roots

- simple root is represented by 0, if short colour the circle
- 90 degrees between, don't connect. $120, 135, 150$ are 1 line, 2 line, 3 line connections respectively
- no loops

- no circle can have 3x lines emanating from it } quite powerful
- removing a circle + restringing \rightarrow valid diagram
- if there is a 263 line connection between 2 circles all circles on one side are short, and on the other side they are long.
- cutting a line gives diagrams of 2 sub-algebras

see p 50 for some (actually odd) diagrams

Casimir operators

we define a casimir operator as one which commutes with all the elements of the algebra. we can always define such an operator
 firstly, define the killing metric:

$$g_{\mu\nu} = f_{\mu a}^\rho \cdot f_{\nu\rho}^\alpha \quad (\text{implied summation})$$

we define its inverse as $g^{\mu\nu}$, and can thus construct the casimir oper.

$$C = g^{\mu\nu} X_\mu X_\nu \quad (\text{again, summation implied})$$

which can be proven that it commutes. (proof, see p 52 though won't be asked)

9.3.2: projective representations and lie algebras

the equivalent of a projective repr. we will discuss at the level of the lie algebra. for a phase ϕ we introduce a new generator

$$X_\phi = \sum_i \frac{\partial f_a(x, a_i, \phi)}{\partial \phi} \cdot \frac{\partial}{\partial x_i}$$

which commutes with everything (the other generators) but results in non-trivial structure factors f_{ab}^ϕ .

the specific form of these f_{ab}^ϕ is specified through the jacobi identity
 there is an example of this being done on p52